# Defect Structures in the Brannerite-Type Vanadates. IV. The Crystal Structure of $\mathbf{M n}_{1-x} \boldsymbol{\emptyset}_{x} \mathbf{V}_{2-2 x} \mathbf{M o}_{2 x} \mathbf{O}_{6} ; \boldsymbol{x}=\mathbf{0 . 5 3}$ 

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#### Abstract

The crystal structure of $\mathrm{Mn}_{1-x} \emptyset_{x} \mathrm{~V}_{2-2 x} \mathrm{Mo}_{2 x} \mathrm{O}_{6} ; x=0.53$ of the brannerite type has been refined to $R=$ 0.029 . The space group is $C 2$. The bond length-bond strength calculations indicate a short-range ordering of $\mathrm{Mo}^{6+}$ and vacancies, resulting in the formation of $\mathrm{Mo}-\mathrm{O}-\emptyset-\mathrm{O}-\mathrm{Mo}$ clusters which are distributed at random in the host vanadate structure.


## Introduction

In a recent paper it was shown that one may substitute Mo for V in the $\mathrm{MnV}_{2} \mathrm{O}_{6}$ structure to produce phases of the formula $\mathrm{Mn}_{1-x} \emptyset_{x} \mathrm{~V}_{2-2 x} \mathrm{Mo}_{2 x} \mathrm{O}_{6}$ in which $\emptyset$ represents a vacancy at the $\mathbf{M n}^{2+}$ site and $\mathbf{M o}^{6+}$ replaces the $\mathrm{V}^{5+}$ ions. Both vacancies and Mo are distributed at random over Mn and V sites respectively (1).

In this paper we have undertaken a structural refinement of this compound, with $x$ approximately equal to 0.5 , in an attempt to discover how the introduction of defects into $\mathrm{MnV}_{2} \mathrm{O}_{6}$ affects its structure. Some structural rearrangement might well be expected in the vicinity of both vacancies and Mo cations, as well as local ordering of the Mo ions and vacancies in an effort to com-

[^0]pensate for the excess positive charge introduced by $\mathbf{M o}^{6+}-\mathrm{V}^{5+}$ replacement. Finally, the structure of the crystal surface will be discussed in anticipation of the results proving useful in the study of the relationship between structural and catalytic properties of the system.

## Experimental

Single crystals were grown from a melt of 0.3 mole\% of prereacted $\mathrm{MnV}_{2} \mathrm{O}_{6}$ and 1.4 mole $\% \mathrm{MoO}_{3}$, making use of the sublimation of $\mathrm{MoO}_{3}$ at higher temperatures (see the phase diagram established for $\mathrm{MnV}_{2} \mathrm{O}_{6}-$ $\mathrm{MoO}_{3}$ system in (1)). The crystallization was performed in an open quartz tube by heating the mixture at $600^{\circ} \mathrm{C}$ for 20 hr . After cooling, a large number of dark brown, plate-shaped crystals were obtained. The Mn : V : Mo atomic ratio in the crystals, as determined by atomic absorption spectros-
copy, was $0.42: 0.96: 1.00$, giving an average $x$ value of 0.53 .

Attempts at grinding a crystal to a sphere for the X-ray experiment were unsuccessful, and a plate-shaped morphology was obtained throughout. The largest face corresponded to (201).

A crystal of dimensions $0.20 \times 0.13 \times$ 0.05 mm was used to collect intensity data up to $\Theta=30^{\circ}$ using an Enraf-Nonius CAD 4 diffractometer and graphite-monochromated $\mathrm{Mo} K \alpha$ radiation in the $\theta / 2 \theta$ scan mode and a scan range given by ( $0.80+$ $0.42 \tan \theta)^{\circ}$. Unit cell parameters (Table I) were obtained by least-squares refinement of the setting angles for 15 reflections used in the autoindexing procedure. It has been established from a least-square fit in ( $l$ ) that Vegard's law is obeyed according to:

$$
\begin{aligned}
& a=9.310+0.1738 x \\
& b=3.537+0.1874 x
\end{aligned}
$$

in the range from 0.00 to 0.40 . This gives an average value of $x=0.57$ for our single crystal.

The fluctuation of intensity for two standard reflections was less than $3 \%$ throughout. The intensities were corrected for geometrical factors but not for absorption. Of the 943 reflections measured 712 unique values were obtained by the sort-merge procedure used. One hundred sixty-six reflections with $\left|F_{0}\right| \leq 4 \sigma\left(F_{0}\right)$ were treated as unobserved.

TABLE I
Crystal Data

| $\mathrm{Mn}_{1}-\emptyset_{x} \mathrm{~V}_{2-2 x} \mathrm{Mo}_{2 x} \mathrm{O}_{8} ;$ | $x=0.53(1)$ |
| :--- | :--- |
| monoclinic | $D_{x}=4.185 \mathrm{Mg} \mathrm{m}^{-3}$ |
| $a=9.412(5) \AA$ | $Z=2$ |
| $b=3.64(1) \AA$ | $F(000)=248$ |
| $c=6.767(3) \AA$ | $\mu(\mathrm{MoK})=6.179 \mathrm{~mm}^{-1}$ |
| $\beta=112.00(4)^{\circ}$ | Systematic absences $h k l ;$ |
| $U=215.2 \AA^{3}$ | $h+k=2 n+1$ |
|  | Space group $C 2\left(C_{2}^{3}\right)$ |

It was found that reflections with $h+k$ odd were absent suggesting that the crystal was $C$ centered. No other systematic absences were observed. Three possible space groups were considered, $C 2 / m, C 2$, and $C m$. Refinement in $C 2 / m$ gave a large $R$ factor ( $>8 \%$ ) and it was observed that on Fourier maps the oxygen peaks appeared doubled. Since there was no evidence, either from powder diffraction or single-crystal analysis, for a superstructure, we could safely reject $\mathrm{C} 2 / \mathrm{m}$. Cm was rejected as no suitable structural model could be found in this space group, nor is this space group known in any similar compounds. We therefore refined the structure in $C 2$. The metal atoms were initially assumed, by analogy with the similar brannerite-type vanadates, to be in the positions:

| Mn | $2 a$ | 0 | 0 | 0 |
| :--- | :--- | :--- | :--- | :--- |
| V | $4 c$ | 0.3 | 0.5 | 0.35 |
| Mo | $4 c$ | 0.3 | 0.5 | 0.35 |

TABLE II
Final Fractional Coordinates and Thermal Parameters ( $\times 10^{4}$ )

| Atom | Position | $x$ | $y$ | $z$ | $U_{11}$ | $U_{22}$ | $U_{33}$ | $U_{23}$ | $U_{13}$ | $U_{12}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Mn | $2 a$ | 0 | 0 | 0 | $142(5)$ | $191(6)$ | $152(5)$ | 0 | $41(4)$ | 0 |
| Mo, V | $4 c$ | $3139.2(4)$ | $5018(8)$ | $3502(1)$ | $95(1)$ | $71(1)$ | $114(1)$ | $8(4)$ | $58(1)$ | $-2(3)$ |
| $\mathrm{O}(1)$ | $4 c$ | $1705(3)$ | $4444(9)$ | $1170(5)$ | $114(8)$ | $210(11)$ | $190(8)$ | $-26(8)$ | $38(7)$ | $-27(8)$ |
| $O(2)$ | $4 c$ | $4749(3)$ | $5237(11)$ | $2950(4)$ | $115(7)$ | $115(9)$ | $255(8)$ | $-12(10)$ | $101(6)$ | $-11(10)$ |
| $O(3)$ | $4 c$ | $3084(3)$ | $42(11)$ | $4375(4)$ | $187(7)$ | $70(7)$ | $243(8)$ | $-34(11)$ | $58(6)$ | $45(11)$ |

Note. Estimated standard deviations in parentheses; $T=\exp \left[-2 \pi^{2}\left(U_{11} a^{* 2} h^{2}+U_{22} b^{* 2} k^{2}+U_{33} c^{* 2 l^{2}}+\right.\right.$ $\left.\left.2 U_{23} b^{*} c^{*} k l+2 U_{13} a^{*} c^{*} h l+2 U_{12} a^{*} b^{*} h k\right)\right]$.

A difference Fourier map gave corrected values for these coordinates and revealed, in addition, approximate positions for the oxygen atoms. Least-squares refinement was carried out by the SHELX-76 program (2). The coordinates of the $\mathrm{V}, \mathrm{Mo}$, and O atoms and the anisotropic temperature factors of $\mathrm{Mn}, \mathrm{V}$, Mo and O were refined, together with the site-occupancy factors of each atom, constrained to fit the chemical formula. The final values of the agreement factors, defined as $R=\Sigma|\Delta F| / \Sigma\left|F_{0}\right|$ and $R_{\mathrm{G}}$ $=\left[\Sigma w(\Delta F)^{2} / \Sigma w F^{2}\right]^{1 / 2}$, were $R=0.029$ and $R_{\mathrm{G}}=0.038$. The weighting scheme was $w=$ $\left[\sigma\left(F_{0}\right)\right]^{-1}$. The final results are listed in Table II. ${ }^{1}$ From the refined site-occupancy factors we found that $x=0.53$, a value in good agreement with the results of atomic absorption measurements. The larger value of $x=0.57$ found from Vegard's law is probably less accurate, since the linear expansion of the lattice is known only for $x$ smaller than 0.40 .

## Discussion

The present compound was found to be isomorphous with other vanadates of the brannerite type. Some of them, such as $\mathrm{MgV}_{2} \mathrm{O}_{6}$ (3), $\mathrm{CdV}_{2} \mathrm{O}_{6}$ (4), and $\mathrm{NaVMoO}_{6}$ (5) have been refined in space group $C 2 / m$; $\mathrm{ZnV}_{2} \mathrm{O}_{6}$ (4) on the other hand was refined in space group C2, as in our case.

The bond lengths and angles are listed in Table III. Figure 1 shows the projection of the structure onto ( 010 ) and a stereo view of part of the $\mathrm{Mo} / \mathrm{V}-\mathrm{O}$ chain along [010].

[^1]No evidence for ordering of Mo or vacancies is found in the refinement. The coordination of the oxygen atoms around $M(M=$ Mo or $V$ ) is difficult to assign and it may be described as four-, five-, or sixfold, since two $M-O$ bonds are much longer than the other four. If the coordination is regarded as sixfold the $M \mathrm{O}_{6}$ octahedra exhibit a strong characteristic distortion. They contain two very short $M-O$ bonds of 1.661 and $1.693 \AA$ to terminal $O(1)$ and quasiterminal

TABLE III
Bond Lengths ( $\AA$ ) and Angles $\left({ }^{\circ}\right)$
(a) Environment of $M=\mathrm{Mo}, \mathrm{V}$

| M | -O(1) | 1.661(3) | $\mathrm{O}(2)-M-\mathrm{O}(1)$ | 105.8(1) |
| :---: | :---: | :---: | :---: | :---: |
| $M$ | -O(2) | 1.693(3) | $\mathrm{O}(3)-M-\mathrm{O}(1)$ | 93.9(1) |
| M | -O(3) | $1.913(5)$ | $\mathrm{O}(3)^{1}-M-\mathrm{O}(1)$ | 107.8(2) |
| M | -O(3) | $1.930(5)$ | $\mathrm{O}(3)-M-\mathrm{O}(2)$ | 104.1(2) |
| $M$ | -O(3) ${ }^{\text {Hii }}$ | 2.151(5) | $\mathrm{O}(3)^{\mathrm{ii}}-M-\mathrm{O}(2)$ | 98.7(1) |
|  |  |  | $\mathrm{O}(3)^{12}-M-\mathrm{O}(3)$ | 142.9(2) |
| M | -O(2) | 2.479(4) | $\mathrm{O}(1)-\mathrm{M}-\mathrm{O}(3)^{\mathrm{Hi}}$ | 100.8(1) |
|  |  |  | $\mathrm{O}(2)-\mathrm{M}-\mathrm{O}(3)^{111}$ | 153.4(2) |
| O(1) | - $\mathrm{O}(2)$ | 2.675(4) | $\mathrm{O}(3)-M-\mathrm{O}(3)^{\text {iii }}$ | 73.4(1) |
| O(1) | - $\mathrm{O}(3)$ | 2.618(4) | $\mathrm{O}(3)^{\mathrm{Hi}}-\mathrm{M}-\mathrm{O}(3)^{\mathrm{jii}}$ | 73.1(1) |
| O(1) | -O(3) ${ }^{\text {ii }}$ | 2.905 | $\mathrm{O}(1)-M-\mathrm{O}(2)^{1}$ | 162.8(2) |
| $\mathrm{O}(2)$ | - $\mathrm{O}(3)$ | 2.754 | $\mathrm{O}(2)-M-\mathrm{O}(2)^{\text {i }}$ | 75.7(1) |
| $\mathrm{O}(2)$ | - $\mathrm{O}(3)^{1 \mathrm{i}}$ | 2.753 | $\mathrm{O}(3)-M-\mathrm{O}(2)^{\text {i }}$ | 80.6(1) |
| $\mathrm{O}(3)$ | - $\mathrm{O}(3)^{\text {ii }}$ | 3.643 | $\mathrm{O}(3)^{\text {ii }}-\mathrm{M}-\mathrm{O}(2)^{\text {r }}$ | $77.1(1)$ |
| O(1) | ) $-\mathrm{O}(3)^{\text {ii1 }}$ | 2.954 | $\mathrm{O}(3)^{\text {iii }}-\mathrm{M}-\mathrm{O}(2)^{\mu}$ | 76.3(1) |

$O(2)-O(3)^{\text {iii }} 3.742$
$\mathrm{O}(3)-\mathrm{O}(3)^{\text {jii }} 2.436(3)$
$O(3)^{1 i}-O(3)^{\text {iii }} 2.437(3)$
$\mathrm{O}(1)-\mathrm{O}(2)^{i} 4.095$
$\mathrm{O}(2)-\mathrm{O}(2)^{i}$ 2.634(6)
$\mathrm{O}(3)-\mathrm{O}(2)^{i} \quad 2.873$
$\mathrm{O}(3)^{1 \mathrm{i}}-\mathrm{O}(2)^{\mathrm{i}} \quad 2.781$
$\mathrm{O}(3)^{\mathrm{iii}} \mathbf{O}(2)^{1} 2.919$
(b) Environment of Mn

| $\mathrm{Mn}-\mathrm{O}(1)$ | $2.207(3)$ | $\mathrm{O}(1)-\mathrm{Mn}-\mathrm{O}(1)^{\mathrm{v}}$ | $100.7(2)$ |
| :--- | :--- | :--- | ---: |
| $\mathrm{Mn}-\mathrm{O}(2)^{\mathrm{V}}$ | $2.098(3)$ | $\mathrm{O}(1)-\mathrm{Mn}-\mathrm{O}(2)^{\mathrm{Vi}}$ | $89.2(1)$ |
| $\mathrm{Mn}-\mathrm{O}(1)^{\mathrm{v}}$ | $2.519(3)$ | $\mathrm{O}(1)-\mathrm{Mn}-\mathrm{O}(2)^{\mathrm{iv}}$ | $87.4(1)$ |
|  |  | $\mathrm{O}(1)^{\mathrm{v}}-\mathrm{Mn}-\mathrm{O}(2)^{\mathrm{iv}}$ | $91.1(2)$ |
|  | $\mathrm{O}(1)^{\mathrm{v}}-\mathrm{Mn}-\mathrm{O}(2)^{\mathrm{vi}}$ | $92.7(2)$ |  |
|  | $\mathrm{O}(1)-\mathrm{Mn}-\mathrm{O}(1)^{\text {vii }}$ | $85.6(1)$ |  |
|  |  | $\mathrm{O}(1)^{\mathrm{v}}-\mathrm{Mn}-\mathrm{O}(1)^{\mathrm{viil}}$ | $73.1(1)$ |

[^2]

Fig. 1. (a) Projection of the $\mathrm{Mn}_{1-x} \emptyset_{x} \mathrm{~V}_{2-2 x} \mathrm{Mo}_{2 x} \mathrm{O}_{6}$ structure onto the (010) plane. The dashed line from $\mathrm{Mo} / \mathrm{V}$ to $\mathrm{O}(2)^{4}$ corresponds to $2.479 \AA$, and because it is so much longer than the other nearestneighbor distances, it is not considered to be within the first coordination sphere. (b) Stereo view of part of Mo/V-O chain along [010].

O (2) atoms, two medium $\mathrm{M}-\mathrm{O}$ bonds of $1.913 \AA$ and $1.930 \AA$ to the bridge-forming oxygens $O(3)$ and $O(3)^{i i}$ and two very long $M-\mathrm{O}$ distances of 2.151 and $2.479 \AA$. The last two bonds result from the small elec-tron-donating character of the bridging $O(3)$ and terminal $O$ (2) atoms, respectively. $O$ (2) forms, in this way, a very asymmetric $M$ -$O(2)-M$ bridge, whereas $O(3)$ acquires coordination number 3.
Table IV shows the most important metal-oxygen interatomic distances in the structure, and the corresponding bond strength values, calculated from the equation:

$$
\begin{equation*}
s=\left(R / R_{1}\right)^{-N}, \tag{1}
\end{equation*}
$$

derived by Brown and Shannon (7), where $s$ is the strength and $R$ is the length of an appropriate bond. The functional relationships which relate bond strength to bond length have been derived by fitting the constants $R_{1}$ and $N$ so that the sums of the bond strength around the cations equal as nearly as possible the valency in a large variety of oxide structures ( 43 for $\mathrm{V}, 50$ for Mo). Brown and Shannon showed that the bond-strength sums around any atom normally lie within $5 \%$ of the valency value. The calculation of the bond-strength sums
provides thus a convenient method of testing the plausibility of a proposed crystal structure. In our case, the discrepancies are 6 and $5 \%$ for Mo and V respectively, and moreover they are opposite in sign. This might indicate that the refined $M-\mathrm{O}$ distances reflect the weighted mean between longer $\mathrm{Mo}-\mathrm{O}$ and shorter $\mathrm{V}-\mathrm{O}$ distances.
The bond lengths and bond strengths determined for the well-established bran-nerite-type structure of pure $\mathrm{MgV}_{2} \mathrm{O}_{6}$ (3) are given for comparison in Table IV. An inspection of this table shows a striking increase in bond strength of the two shortest $M-\mathrm{O}(1)$ and $M-\mathrm{O}(2)$ bonds when $\mathrm{Mo}^{6+}$ is substituted for $\mathrm{V}^{5+}$. The bond strengths of the other four bonds remain practically unchanged. As to $\mathrm{VO}_{6}$ octahedra, the bonding scheme of $\mathrm{V}^{5+}$ ions differs insignificantly from that in the pure vanadates $\left(\mathrm{MgV}_{2} \mathrm{O}_{6}\right)$. The observed effects can be accounted for by assuming that there occurs a short-range ordering of $\mathbf{M o}^{6+}$ and vacancies resulting in the formation of $\mathrm{Mo}-$ $\mathrm{O}-\varnothing$-O-Mo clusters, which are distributed at random in the host vanadate structure. The effect may be illustrated by the bond-strength sums calculated around $\mathrm{O}(1)$ and $\mathrm{O}^{(2)}$ (Table V ). All the sums are about 2 if we assume that $\mathrm{Mo}^{6+}$ involves a va-

TABLE IV
Bond Strength Sums in Some Brannerite-Type $A B_{2} \mathrm{O}_{6}$ Compounds ${ }^{a}$

| $\mathrm{Mn}_{1-x} \emptyset_{x} \mathrm{~V}_{\mathbf{2}-\mathrm{8x}} \mathrm{Mo}_{2 x} \mathrm{O}_{8} ; x=0.53$ |  |  | $\mathrm{MgV}_{2} \mathrm{O}_{\mathbf{B}}$ (Ref. (3)) |  | $\mathrm{NaVMoO}_{6}$ (Ref. (5)) |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{gathered} M-\mathrm{O} \\ {[\AA]} \end{gathered}$ | $s_{\text {M } 0^{\text {a }}}$ | $s^{\text {v }}{ }^{\text {s+ }}$ | $\begin{aligned} & \text { V-O } \\ & {[\AA]} \end{aligned}$ | $s^{\text {v }}$ st | $\begin{gathered} M-\mathrm{O} \\ {[\AA]} \end{gathered}$ | $S_{\text {moat }}$ |
| 1.661 | 2.115 | 1.469 | 1.671 | 1.424 | 1.63 | 2.37 |
| 1.693 | 1.887 | 1.332 | 1.666 | 1.446 | 1.73 | 1.66 |
| 1.913 | 0.907 | 0.715 | 1.853 | 0.841 | 1.92 | 0.89 |
| 1.930 | 0.860 | 0.683 | 1.853 | 0.841 | 1.92 | 0.89 |
| 2.151 | 0.449 | 0.393 | 2.111 | 0.432 | 2.15 | 0.45 |
| 2.479 | 0.191 | 0.191 | 2.671 | 0.130 | 2.47 | 0.20 |
| 玉 | 6.409 | 4.783 |  | 5.114 |  | 6.46 |
| $\begin{gathered} \mathrm{Mn}-\mathrm{O} \\ {[\AA]} \end{gathered}$ | $s_{\text {Mn }}+2$ |  | $\begin{gathered} \mathrm{Mg}-\mathrm{O} \\ {[\AA]} \end{gathered}$ | $s_{\text {Mg }^{2+}}$ | $\begin{gathered} \mathrm{Na}-\mathrm{O} \\ {[\AA]} \end{gathered}$ | $s_{\mathrm{Na}^{+}}$ |
| $2.098(2 \times)$ | 0.421 |  | 2.024(2x) | 0.429 | 2.29(2x) | 0.228 |
| $2.207(2 x)$ | 0.317 |  | 2.199(2x) | 0.271 | 2.44( $2 \times$ ) | 0.173 |
| $2.519(2 \times)$ | 0.151 |  | 2.199(2x) | 0.271 | 2.44(2x) | 0.173 |
| $\Sigma_{s}$ | 1.778 |  |  | 1.942 |  | 1.148 |

${ }^{a}$ Bond strength equations from ( 6 ) $; s=(R / 1.882)^{-8.0}$ for $\mathrm{Mo}^{6+}, s=(R / 1.791)^{-5.1}$ for $\mathrm{V}^{5+}, s=(R / 1.798)^{-5.6}$ for $\mathrm{Mn}^{+2}, s=(R / 1.622)^{-4.29}$ for $\mathrm{Mg}^{2+}$ and $\mathrm{Na}^{+}$.
cancy whereas $\mathrm{V}^{5+}$ involves a $\mathrm{Mn}^{2+}$ cation bridges, compared with the $\mathrm{MgV}_{2} \mathrm{O}_{6}$ strucin a neighboring $M^{2+}$ site.

The lengthening of the $M-O(3)$ and $M-$ $\mathrm{O}(3)^{11}$ distances in the symmetrical $M-O-M$ ture, is apparently due to changes in ionic size (the ionic radii of $\mathrm{Mo}^{6+}$ and $\mathrm{V}^{5+}$ are 0.60 and $0.54 \AA$ respectively (8)). This

TABLE V
Oxygen Bond-Strength Sums in the Postulated
Model of $\mathrm{Mn}_{1-x} \partial_{x} \mathrm{~V}_{2-2 x} \mathrm{Mo}_{2 x} \mathrm{O}_{6}{ }^{a}$

| Postulated $O(1)$ and $O(2)$ environments | $\Sigma_{s}$ |
| :--- | :--- |

[^3]statement is confirmed by the increase of the unit cell parameter $b$ parallel to the symmetric $M-O(3)-M$ bridges, in accordance with Vegard's law, as found in the earlier work (1).

The $M-O(2)$ bond in the solid solution acquires a more pronounced bridging character compared with that in pure vanadates. The two very short bonds to the terminal oxygens, which are of almost equal length in pure vanadates, become nonequivalent. The increase in the $M-O(2)$ distance brings about a shortening of the $M-O(2)^{i}$ distance to a value of $2.48 \AA$ in order to maintain the bond strength around the $M$ ions.

The distances found for the investigated compound fall fairly close to those determined for the isomorphous $\mathrm{NaVMoO}_{6}$ structure (5) (see table IV); here $\mathrm{Na}^{+}$ions occupy the $M^{2+}$ positions, whereas $\mathrm{V}^{5+}$ and $\mathrm{Mo}^{6+}$ are distributed at random over the $M$ positions. The only difference is the shortening of the $M-\mathrm{O}(1)$ bond to a value of 1.63 $\AA$, accompanied by the lengthening of the $M-O(2)$ bond to $1.73 \AA$.

If we consider only five bonds involved in the closer coordination of $M$, with the exclusion of the very long $M-O(2)^{i}$ bond, the dominating element in the structure is an infinite double chain of edge-linked distorted $M \mathrm{O}_{5}$ tetragonal pyramids parallel to the $b$ axis. The distortion of the fivefold coordination around the V and Mo atoms is due to their displacements in opposite directions, in order to minimize electrostatic repulsions, and the partly covalent character of Mo-O and V-O bonds. The double chains are linked to each other by $\mathrm{Mn}^{2+}$ atoms ( $\mathrm{Mn}-\mathrm{O}(2)$ bonds) forming sheets parallel to the (201) plane. The sheets are
bound together by weak $\mathrm{Mn}-\mathrm{O}(1)$ and $M-$ $O(2)^{i}$ bonds. The structure should show good cleavage on (201), ${ }^{2}$ and this was confirmed in this work by the observed plate-like morphology of the investigated single crystals. The preferential cleavage on (201) exposes the $M=0$ groups on the surface. The tendency of molybdenum to produce isolated $\mathrm{Mo}=\mathrm{O}$ covalent bonds should bring with it the scope for increased concentration of Mo in the surface beyond the average concentration in the bulk of the structure.

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[^1]:    ${ }^{1}$ See NAPS document No. 03838 for 3 pages of supplementary material. Order from ASIS/NAPS c/o Microfiche Publications, P.O. Box 3513, Grand Central Station, New York, N.Y. 10017. Remit in advance in U.S. funds only for each NAPS Accession Number. Institutions may use purchase orders when ordering; however, there is a billing charge for this service. Make checks payable to Microfiche Publications. Photocopies are $\$ 5.00$. Microfiche are $\$ 3.00$. Outside of the U.S. and Canada, postage is $\$ 3.00$ for a photocopy or $\$ 1.50$ for a fiche.

[^2]:    Note. (i) $1-x, y, 1-z$; (ii) $x, 1+y, z$; (iii) $\frac{1}{2}-x$, $\frac{1}{2}+y, 1-z$; (iv) $-\frac{1}{2}+x,-\frac{1}{2}+y, z ;$ (v) $x,-1+y, z$; (vi) $\frac{1}{2}-x,-\frac{1}{2}+y,-z$; (vii) $-x,+y,-z$; (viii) $-x$, $-1+y,-z$.

[^3]:    ${ }^{a}$ Bond-strength equations as in Table IV

[^4]:    ${ }^{2}$ The preferential cleavage on (201) was kindly pointed out to us by Dr. J. Ziółkowski on the basis of preferred orientation effects in X-ray powder diffractometer samples.

